

Reversible Water-Induced Phase Changes of Cobalt Oxide Nanoparticles

*Dennis Hein^{a,b}, Garlef Wartner^{a,b}, Arno Bergmann^c, Miguel Bernal^{c,d,†}, Beatriz Roldan Cuenya^c,
and Robert Seidel^{a,b,*}*

^a Helmholtz-Zentrum Berlin für Materialien und Energie (HZB), Albert-Einstein-Straße 15,
12489 Berlin, Germany

^b Humboldt-Universität zu Berlin, Institut für Chemie, Brook-Taylor-Str. 2, 12489 Berlin,
Germany

^c Fritz-Haber-Institut der Max-Planck-Gesellschaft, Faradayweg 4-6, 14195, Berlin, Germany

^d Department of Physics, Ruhr-University Bochum, 44780 Bochum, Germany

* corresponding author

† **Present Address**

Faculté des Sciences, Chemistry of Surfaces, Interfaces and Nanomaterials (ChemSIN), Université
Libre de Bruxelles, Avenue Franklin Roosevelt 50, 1050 Bruxelles, Belgium.

ABSTRACT. Cobalt oxides have been identified as highly active catalysts for the electrochemical
water splitting and oxygen evolution reaction. Using near-ambient pressure resonant photoelectron

spectroscopy, we studied changes in the metal-oxygen coordination of size-selected core-shell CoO_x nanoparticles induced by liquid water. In dry conditions, the nanoparticles exhibit an octahedrally coordinated Co^{2+} core and a tetrahedrally coordinated Co^{2+} shell. In the presence of liquid water, we observe a reversible phase change of the nanoparticle shell into octahedrally coordinated Co^{2+} as well as partially oxidized octahedrally coordinated Co^{3+} . This is in contrast to previous findings, suggesting an irreversible phase change of tetrahedrally coordinated Co^{2+} after the oxygen evolution reaction conditioning. Our results demonstrate the appearance of water-induced structural changes different from voltage-induced changes and help us to understand the atomic scale interaction of CoO_x nanoparticles with water in electrochemical processes.

KEYWORDS. near-ambient pressure XPS, solid-liquid interface, cobalt oxide, reversible phase change, core-shell nanoparticles

The electrochemical water splitting process consists of two half reactions, namely the hydrogen evolution reaction (HER) and oxygen evolution reaction (OER) for which commonly expensive and rare noble elements like Pt and Ir/Ru based catalysts are used, respectively, hindering the commercial application.^{1,2} For this reason, extensive research efforts are made to find abundant, active and stable electrocatalysts, with 3d transition metal oxides being a promising material class.^{1,2} Key catalytic performance parameters such as the efficiency and stability, depend on molecular processes determined by the electronic-structure at the solid-liquid interface between the electrocatalyst and the electrolyte. An in-depth insight into such an interface and its evolution under different chemical environments and under potential control can be now achieved thanks to modern spectroscopic tools such as near-ambient pressure photoelectron spectroscopy (NAP-XPS).

Due to its abundance and catalytical activity we focus on cobalt oxide nanoparticles (NPs).^{1,3-11} The investigation of different cobalt oxide NPs revealed that they have similar catalytic activities for the oxygen evolution reaction (OER) independent of their oxidation state and coordination, with Co_3O_4 being the most stable against corrosion.¹ However, it has been shown that the near-surface oxygen chemistry of $\text{CoO}_x(\text{OH})_y$ and redox electrochemistry are correlated to the catalytic activity, showing the importance of reducible μ_2 -OH bridged Co^{3+} sites.¹²

Bulk cobalt oxide is thermodynamically stable at room temperature in two forms: 1) rocksalt CoO , which consists of octahedrally (O_h) coordinated Co^{2+} , and 2) spinel Co_3O_4 , consisting of the tetrahedral (T_d) coordinated Co^{2+} A site and the O_h Co^{3+} B site.¹³⁻¹⁵ Selectively replacing a phase of spinel Co_3O_4 with inactive Zn^{2+} or Al^{3+} revealed that under OER conditions O_h Co^{3+} forms a stable bond with -OH groups, limiting its catalytical activity, whereas Co^{2+} forms CoOOH ^{16,17}, as predicted by theory¹⁸. This transformation of T_d Co^{2+} towards CoOOH in spinel Co_3O_4 under applied oxidizing bias is a **reversible** process, induced by its oxygen affinity, which is a result of

the ability to release electrons under applied potential.^{16,17,19,20} Near-ambient pressure XPS of Co_3O_4 demonstrated this transformation during OER and its immutability without applied oxidizing bias.²¹ In contrast, $\text{O}_h \text{Co}^{2+}$, the more active site for oxygen reduction reactions (ORR), transforms **irreversibly** into $\text{O}_h \text{Co}^{3+}$ and $\text{T}_d \text{Co}^{2+}$ during ORR.²² Another **irreversible** transformation occurs for T_d and $\text{O}_h \text{Co}^{2+}$ oxides, forming $\text{CoO}_x(\text{OH})_y$ during OER conditioning, contrary to the mentioned **reversible** transformation of Co_3O_4 .¹² These transformations demonstrate the significance of understanding the specific solid-liquid interface that determines the surface structure and surface processes relevant for water splitting applications.

In the case of CoO_x NPs, Papaefthimiou *et al.* suggested that the core consisted of rocksalt CoO ($\text{O}_h \text{Co}^{2+}$) surrounded by a partly wurtzite CoO ($\text{T}_d \text{Co}^{2+}$) shell.¹⁵ Macroscopic wurtzite is not stable under ambient conditions, since the rocksalt structure is energetically favoured (by 0.27 eV per CoO , according to lattice energy calculations²³). However, it can still occur embedded in another crystal²⁴ supporting its structure, or due to surface stress as a metastable phase on the nanoscale length, either as a nanoparticle^{12,14,25,26} or as a thin film on top of epitaxially grown rocksalt CoO .^{27,28} Lukashuk *et al.* demonstrated that the grain size of cobalt oxide influences the flexible adjustment to various reaction environments.²⁹ Additionally, Fischer-Tropsch synthesis studies with CoO_x NPs under reducing and oxidizing atmospheres showed their stability up to a temperature where bulk CoO_x is easily oxidized or reduced in the respective atmospheres, *e.g.* 630 K in O_2 atmosphere³⁰, possibly due to the stabilizing wurtzite shell.^{30,31}

Figure 1 (a) sketches the crystal structure and a simplified 3d orbital occupancy of $\text{O}_h \text{Co}^{2+}$ and $\text{T}_d \text{Co}^{2+}$.¹²⁻¹⁴ Due to a different coordination and crystal field splitting, the order and occupation of the e_g and t_{2g} levels are also different. For the tetrahedrally coordinated $\text{T}_d \text{Co}^{2+}$ case, the lower lying e_g level is fully occupied, whereas the t_{2g} level of the $\text{O}_h \text{Co}^{2+}$ species has a one-electron

vacancy.^{13,32} This vacancy will lead to a pre-peak in the Co 2p X-ray absorption spectra (XAS) of $O_h Co^{2+}$, making it possible to distinguish between both configurations. For comparison, Figure 1 (a) also displays the spinel Co_3O_4 crystal and its electronic structure, which consists of two-thirds $O_h Co^{3+}$.^{12,33} As will be discussed in Figure 2, we focus on the partial electron yield x-ray absorption spectra (PEY-XAS), which are the integral of the specific Auger-electron decay channel signals (see Figure 1) as a function of the photon energy. A detailed description can be found in the methods section.

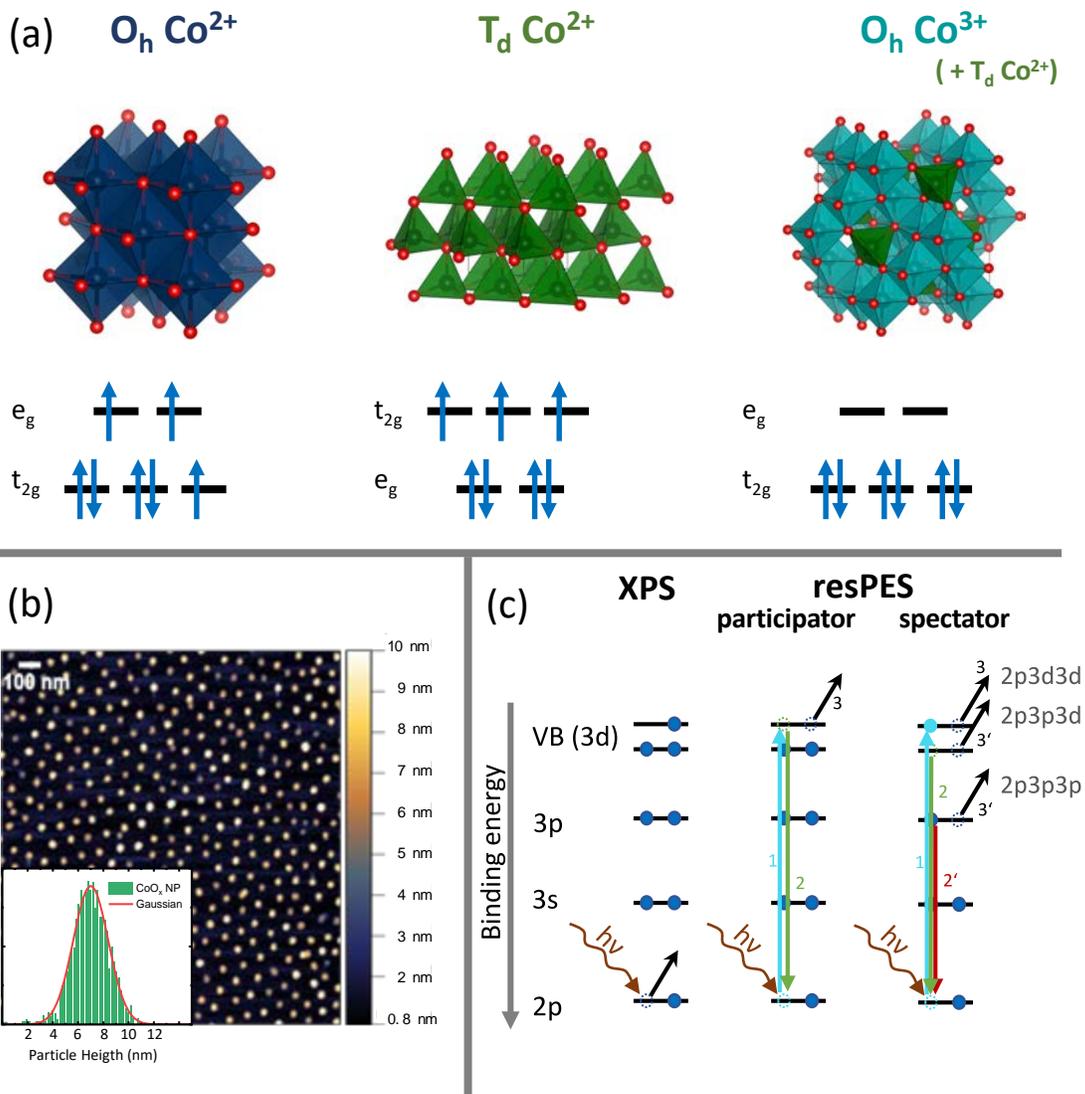


Figure 1. (a) Crystal structure of rocksalt (rs) CoO, wurtzite (wz) CoO, and spinel Co₃O₄. Polyhedrons represent the different coordination structures of the cobalt complexes. Dark blue octahedrons represent O_h Co²⁺, light blue octahedrons in Co₃O₄ are O_h Co³⁺ coordinated, and green tetrahedrons are T_d Co²⁺. The corresponding electron occupation of the e_g and t_{2g} molecular orbitals for each Co-O coordination is shown at the bottom.^{13,34,35} (b) Atomic force microscopy image of the size-selected CoO_x NPs with a height of 7.0 ± 2.0 nm on a SiO_x/Si(111) substrate. (c) Schematic representation of the (normal) XPS process, the participator and the spectator Auger decay. The latter two only occur in resonance.

This observed stability and the suggested core-shell structure of CoO_x NPs raises the question of the specific structural interface in contact with water. The interface properties, especially the near-surface electronic structure, determine the molecular processes during the OER. Naturally, NAP-XPS is the method of choice to experimentally investigate the electronic structure and the element-specific surface composition. For this purpose a few-monolayer thick water film was adsorbed onto the CoO_x NPs solid surface in contact with a SiO_x/Si(111) substrate. Figure 1 (b) displays an AFM image of the used CoO_x NPs on a SiO_x/Si(111) substrate. Several groups explored already in detail the liquid film formation on several metal and metal oxide samples, *e.g.* TiO₂(110), Cu(110), Cu(111)³⁶, Fe₃O₄(001)³⁷ and α-Fe₂O₃(0001)³⁸, as a function of the water vapor pressure, relative humidity and sample temperature, indicating the necessary pressure range to form a liquid H₂O monolayer.

We apply soft X-ray NAP-XPS to size-selected CoO_x NPs, focusing on resonant X-ray photoelectron spectroscopy (resXPS) at the cobalt L₃-edge, to gain insight into the water/CoO_x NP

interface. Here, we report a **reversible** water-induced phase-change between tetrahedrally coordinated Co^{2+} into octahedrally coordinated Co^{2+} and partially oxidized octahedrally coordinated Co^{3+} , occurring **without** an applied oxidizing bias. Additionally, we discuss the resonant enhanced valence-band structure of the catalyst in contact w/ and w/o liquid H_2O .

RESULTS AND DISCUSSION

PEY-XAS under high vacuum and near-ambient pressure conditions (Co L_3 -edge).

Information on the oxidation state and atomic coordination of the CoO_x NPs can be achieved by sweeping the photon energy through the Co L_3 -edge and integrating the resonant Auger decay channels to generate the 2p3p3p and 2p3p3d PEY-XAS, presented in Figure 2 (a) and (b), respectively. The notation, *e.g.*, 2p3p3d describes a 2p hole that is refilled by a 3p electron and the excess energy is transferred to a 3d electron that leaves the system. In all measured PEY-XA spectra we subtracted a linear background and normalized to the absorption maximum at around $E_{\text{ph}} = 778.1\text{eV}$. Both black PEY-XA spectra were measured under high vacuum conditions at 2.5×10^{-5} mbar. For comparison, in Figure 2 (c) we present calculated spectra by Bergmann *et al.*¹² for $T_d \text{Co}^{2+}$, $O_h \text{Co}^{2+}$ and $O_h \text{Co}^{3+}$ species, conducted using atomic positions calculated by density functional theory, which were employed into *ab initio* crystal field multiplet calculations³⁹. The absorption spectra of the CoO_x NPs under high vacuum conditions are a superposition of spectra from $T_d \text{Co}^{2+}$ and $O_h \text{Co}^{2+}$ phases, as seen in comparison with the calculated spectra. One spectral fingerprint for $O_h \text{Co}^{2+}$ is the extra intensity at 777.2 eV (marked with **A** in Figure 2), which reflects the different occupation of the e_g and t_{2g} molecular orbitals (see Figure 1 (c)). Due to the constant electron inelastic mean free path (IMFP, $\sim 1.5 \text{ nm}$ ⁴⁰) and the approx. 7 nm size of our NPs, we obtain mainly signals from the outer NP layers, suggesting the existence of an $O_h \text{Co}^{2+}$ and a $T_d \text{Co}^{2+}$ mixed shell or an $O_h \text{Co}^{2+}$ core - $T_d \text{Co}^{2+}$ shell structure, as proposed by Papaefthimiou *et al.*¹⁵

Under NAP conditions at a 4-mbar water atmosphere (in blue), we obtain an increased intensity for the 2p3p3p and the 2p3p3d PEY-XAS at 777.2 eV (marked with **A** in Figure 2) as well as in the photon energy range between 779.6 eV and 784.0 eV (marked with **B**). The increase at around 780.0 eV photon energy is due to the enhancement of the $O_h Co^{2+}$ signal, similarly to the increase at 777.2 eV. The extra intensity at energies beyond 782 eV is a fingerprint of $O_h Co^{3+}$.^{12,41-44}

We cannot exclude completely the formation of spinel Co_3O_4 or $CoOOH$ at NAP conditions as suggested by the $O_h Co^{3+}$ signal, but we do not observe the expected associated decrease of the $O_h Co^{2+}$ signal at 777.2 eV in the PEY-XAS.¹² The obtained spectral features of adsorbed water in Figure 2 (a) are less pronounced for the 2p3p3d PEY-XAS in Figure 2 (b), which is due to a higher delocalization probability of the d-electrons. During the few femtoseconds Co 2p core lifetime, the core-excited electron in the 3d valence band can delocalize with its surrounding, weakening the Auger decay channels involving a 3d electron. Thus, the PEY-XAS investigations suggest a mixed-shell or core-shell structure of the CoO_x NPs at high vacuum pressure conditions and a water-induced phase change of $T_d Co^{2+}$ to $O_h Co^{2+}$ and $O_h Co^{3+}$.

Notably, the obtained extra intensities at near-ambient pressure go back to its initial state after having reestablished the high vacuum conditions, which indicates a **reversible** behavior. Figure 2 (b) shows the 2p3p3d PEY XAS at high vacuum before (black), during (blue), and after (brown) the investigations at NAP conditions. The 2p3p3d PEY-XA spectra under near-ambient conditions and afterwards at high vacuum conditions were measured for these spectra at the same sample spot.

There are two different transformations known for $T_d Co^{2+}$ in contact with aqueous solutions: 1) an **irreversible** transformation of 'single-phased' $T_d Co^{2+}$ towards 3D cross-linked $Co^{2+/3+}O_x(OH)_y$ by OER conditioning¹² and 2) a **reversible** transformation of the $T_d Co^{2+}$ in Co_3O_4 towards $CoOOH$ under OER condition^{16,17,19,20}. As we investigated the CoO_x at open-circuit conditions, both

changes should not take place. But as $T_d \text{Co}^{2+}$ is thermodynamically not favored, an irreversible phase change similar to the findings during OER conditions is expected. Our results now suggest that the transformation of the $T_d \text{Co}^{2+}$ phase in the presence of water reflects a possible intermediate state towards the further transformation during OER, occurring already at open-circuit conditions.

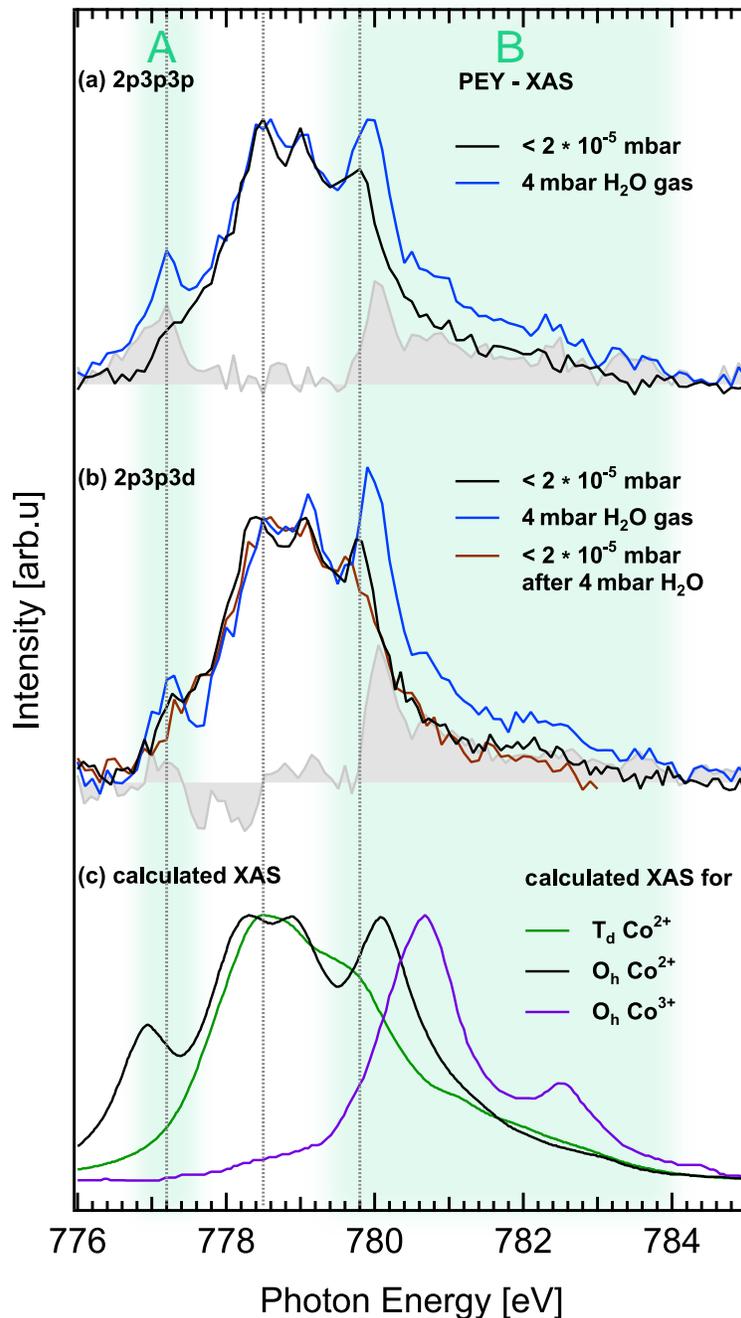


Figure 2. PEY-XAS of (a) 2p3p3p and (b) 2p3p3d of CoO_x NPs at high vacuum (black), at 4 mbar H₂O atmosphere (blue) and again at high vacuum after the 4 mbar H₂O measurement (brown). (c) Calculated XAS for T_d Co²⁺, O_h Co²⁺ and O_h Co³⁺.¹² Raw spectra were corrected with a linear background and normalized to the maximum intensity at E_{ph} = 778.1eV.

Resonant Auger spectra under near-ambient pressure conditions. Figure 3 shows the resonant Auger spectra at the cobalt L_3 -edge under NAP conditions to reveal the subtle differences in electronic-structure during the NP-water interaction. The excitation energies used are $E_{ph} = 772.0$ eV (off-resonant, black curve) and $E_{ph} = 777.2$ eV (on-resonant, blue), enhancing only O_h Co^{2+} contributions of the NP in contact with liquid H_2O (see Figure 2). In the resonant case, the spectator Auger decay causes four spectral increases: (1) the Co 2p3p3p Auger decay at a binding energy larger than 120 eV (electron kinetic energy $E_{kin} < 653$ eV), (2) the Co 2p3s3d Auger decay at a binding energy larger than 100 eV ($E_{kin} < 673$ eV), (3) the Co 2p3p3d Auger decay at binding energies larger than 60 eV ($E_{kin} < 713$ eV) and (4) the Co 2p3d3d Auger decay at binding energies larger than 6 eV ($E_{kin} < 769$ eV). Considering the low surface coverage of the CoO_x NPs (approx. less than 1 %) the spectral increase using a resonant photon energy is significant and extends over a wide binding energy range, because of the large number of possible final states. The 3d-electrons of cobalt contain information of its valence electronic-structural environment, due to hybridization with its surrounding (mainly the lattice oxide), which also leads to a higher delocalization probability of the d-electron. Therefore, Auger decays involving d-electrons sense changes in the electronic environment, but are also weakened, as the core-excited electron in the 3d valence band can delocalize with its surrounding during the Co 2p core lifetime of a few femtoseconds. Consequently, the Co 2p3d3d Auger decay in Figure 3 has the weakest intensity.

Besides the spectral increase of the spectator Auger decays, the participator Auger decay enhances the Co 3p signal as seen in the blue spectrum of Figure 3, with its maximum at $E_B = 63.2$ eV. Similar enhancements of the 3s at around 100 eV⁴⁵ are overlaid by the Si 2p and SiO_x 2p orbitals. We included the calculated off-resonant and resonant 3p spectra by Tanaka *et al.*⁴⁶ for a

cluster composed of O_h Co^{2+} surrounded by six O^{2-} ions for comparison. They used a cluster model with full multiplet splitting for their calculations that matches qualitatively our experimental data.⁴⁶ Although these modeled values do not include a liquid H_2O surrounding, they still predict the obtained data well, except for the spectator Auger spectral region around 70 eV.

As 99% of the sample surface consists of the $SiO_x/Si(111)$ substrate, the spectra in Figure 3 display strong SiO_x features, namely the SiO_x 2p (104.7 eV), the Si 2p_{3/2} (99.4 eV) and the 2p_{1/2} (100.0 eV). At NAP conditions the SiO_x O 2s (27.5 eV) and $SiO_x/Si(111)$ valence band structure are superimposed by H_2O gas and liquid orbitals. Literature values for the H_2O gas peak positions are highlighted in Figure 3.⁴⁷ A detailed analysis of the $SiO_x/Si(111)$ substrate using NAP-XPS can be found in Figure SI-2. Due to the low surface coverage of the CoO_x NPs, off-resonant O 1s and O 2s spectra contain hardly any contributions from either CoO_x NP lattice oxide or from adsorbed H_2O and possible terminated OH-groups on the NP surface.

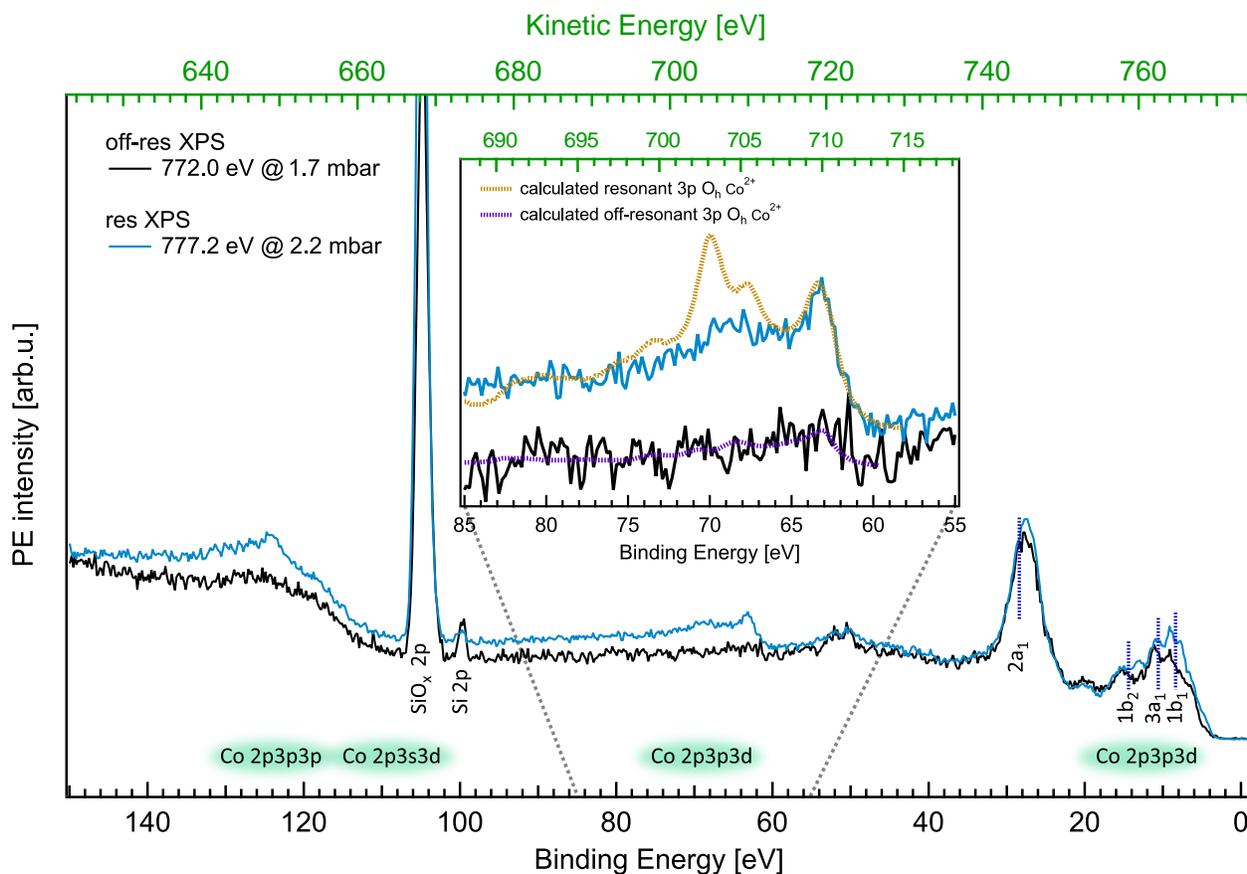


Figure 3. Co 2p \rightarrow 3d on-resonant (777.2 eV, blue curve) and off-resonant (772.0 eV, black curve) photoelectron spectra of CoO_x NP on a SiO_x/Si(111) substrate, measured at \sim 2 mbar pressure. In resonance, the Co 2p_{3p3p}, Co 2p_{3s3d}, Co 2p_{3p3d} and Co 2p_{3d3d} Auger decay are visible. Off- and on-resonant spectra are dominated by SiO_x and Si signals, and also contribution from gas and liquid water are present in the valence band. The binding energies of the water molecular orbitals are marked.⁴⁷ Inset: Zoom-in of Co 2p_{3p3d} Auger decay signal together with calculated on-resonant (yellow) and off-resonant spectra (purple) of O_h Co²⁺.⁴⁶

Resonant-PES at high vacuum vs near-ambient pressure (valence band). Additional studies on the participator Auger decay, measured at defined photon energies, reveal the enhanced 3d spectrum, from which we can derive the different Co-species present in the NP surface. Several pairs of off-and on-resonant spectra were measured, changing to a new sample spot when measuring a new pair at high vacuum and near-ambient pressure conditions. The difference ‘on-resonant spectrum’ minus ‘off-resonant spectrum’ (grey) for several selected photon energies (‘off-resonant’: 772.0 eV and ‘on-resonant’: 777.2 eV, 778.5 eV and 779.8 eV; marked at the integrated PEY-XAS spectra in the insets) are plotted in Figure 4A at high vacuum, and in Figure 4B at near-ambient pressure conditions, as well as in Figure SI-3. Positive values in the difference spectra reveal an enhanced Co 3d signal and 2p3d3d spectator Auger decay. As previously discussed, the valence spectra measured at high vacuum conditions (Figure SI-4) are dominated by the SiO_x/Si(111) contributions. Using the PEY-XAS results, one expects the resonant enhancement at high vacuum to be a superposition of T_d Co²⁺ 3d and O_h Co²⁺ 3d signals. At the lowest resonant photon energy, the calculated XAS predicts a strong O_h Co²⁺ feature which absence indicates a stronger T_d Co²⁺ NP component and PES signal. Additionally, the difference spectra at high vacuum conditions do not resemble the calculated O_h Co²⁺ 3d resonant spectra as shown in Figure 4b, but can be emulated when we include calculated final states for T_d Co²⁺ (green).⁴⁸

The valence spectra measured at NAP conditions (Figure 4b) are dominated by H₂O gas and liquid contributions. For the difference at near-ambient pressure conditions we append the calculated O_h Co²⁺ 3d resonant spectra by Tanaka *et al.* (orange) and calculated final states for O_h Co³⁺ (purple) using ligand-field theory by Jugnet *et al.*^{46,48} The calculated spectra resemble our data, especially for lower binding energies at a photon energy of 777.2 eV. At higher excitation energies a signal contribution around 4.9 eV arises by the ²T_{2g} final state of O_h Co³⁺ (purple stick).

The measured signal in Figure 4b at binding energies between 5-10 eV is lower than the calculated spectra.⁴⁹

These findings confirm that: i) at high vacuum conditions the main NP component is T_d Co²⁺, and ii) a phase transformation from the T_d Co²⁺ NP phase to O_h Co²⁺ and O_h Co³⁺ occurs under NAP conditions, as it is also revealed by the PEY-XA spectra (Figure 2).

X-ray radiation effects and sample contamination. Not yet mentioned is the influence of radiation and contamination affecting the sample. Despite differently noted, we changed between PEY-XAS and resonant XPS measurements regularly the spot to prevent long exposure radiation damage effects. For the PEY-XAS measurements we also lowered the photon flux by one order of magnitude compared to the resXPS measurements. Contamination of our samples by hydrocarbons is inevitable due to an *ex situ* preparation. The existence of oxygen species beside SiO_x is revealed by the O 1s and C 1s XPS (Figure SI-4). Since we use resonant photon energies, the Co 2p→3d PEY-XA spectra are dominated by cobalt signals. To ensure this for the resXPS, we performed off-resonant PES before every on-resonant measurement. Positive values for the difference ‘on-resonant spectrum’ minus ‘off-resonant’ spectrum refer solely to cobalt enhancements. Negative values for the difference are not a result of the resonant PES itself, but of a continuous reduction of the signal, most likely due to a continuously increasing surface contamination of carbon species, CO_xH_y.

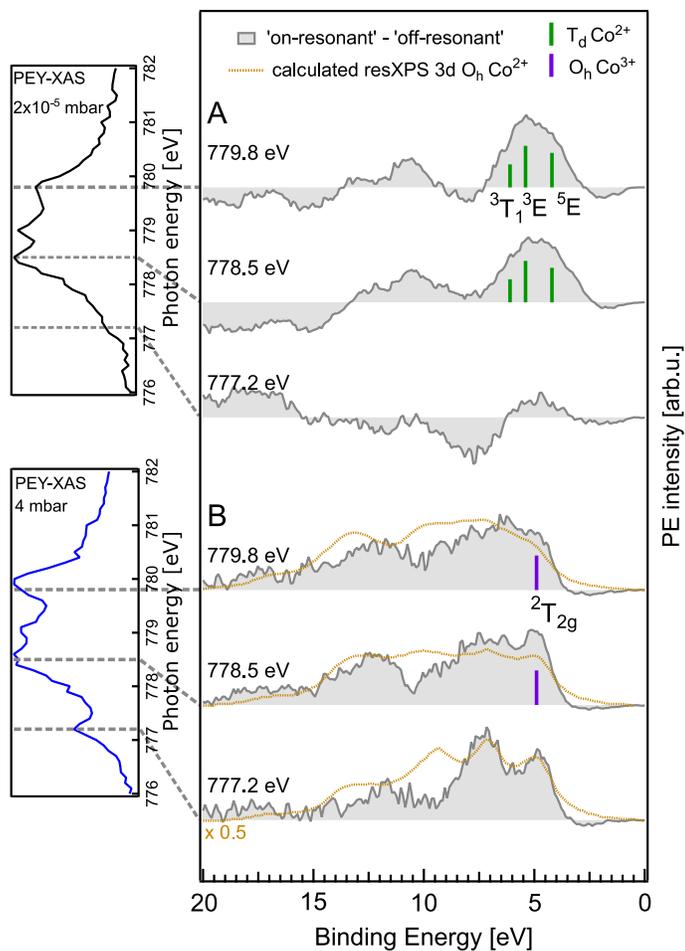


Figure 4. Differences of 'on-resonant' spectrum (777.2 eV, 778.5 eV and 779.8 eV) minus 'off-resonant' spectrum (772.0 eV) XPS taken at (A) high vacuum, and at (B) NAP conditions. Included are calculated 3d resonant XPS of $O_h Co^{2+}$ and calculated 3d final states of $T_d Co^{2+}$ and $O_h Co^{3+}$.^{46,48} The boxes on the left hand side display the 2p3p3p PEY-XAS measurements from Figure 3.

CONCLUSION

We used NAP-XPS to identify the phase changes in size-selected CoO_x NPs on a $\text{SiO}_x/\text{Si}(111)$ substrate. Under high vacuum conditions these NPs consist of an $\text{O}_h \text{Co}^{2+}$ and a $\text{T}_d \text{Co}^{2+}$ phase, with a proposed T_d shell.¹⁵ In a water atmosphere (~ 4 mbar), the composition is altered to an $\text{O}_h \text{Co}^{3+}$ and an $\text{O}_h \text{Co}^{2+}$ phase. When reapplying high vacuum conditions, and after water evaporation, the previous cobalt oxide composition is reestablished. We observe that water changes the shell of the NPs, resulting in the reversible destabilization of the metastable $\text{T}_d \text{Co}^{2+}$ phase and in the formation of $\text{O}_h \text{Co}^{2+}$ and oxidized $\text{O}_h \text{Co}^{3+}$. Since we operate our experiments without applied an oxidizing bias, our measurements reveal an intermediate state. We cannot completely exclude a possible formation of CoOOH , however, we do not observe the necessary accompanied decrease of the $\text{O}_h \text{Co}^{2+}$ signal in the PEY-XA spectra. Using resonant XPS we could also spectroscopically reveal the cobalt oxide valence band of $\text{O}_h \text{Co}^{2+} + \text{T}_d \text{Co}^{2+}$ at high vacuum and of $\text{O}_h \text{Co}^{3+} + \text{O}_h \text{Co}^{2+}$ at near-ambient pressure.

The next logical step would be to carry out the resonant XPS experiments under operando conditions. This requires a different approach, *e.g.* the use of special micro-electrochemical cells with graphene-covered holey-membranes.^{50,51} The few monolayer of graphene act as the working electrode and the carrier for the deposited CoO_x nanoparticles. These experiments are currently planned.

Overall, we showed that a water-induced phase transformation affects the CoO_x NP catalyst differently than voltage-induced changes in an aqueous electrolyte, which is relevant for our understanding of the interaction of Co ions with water at the atomic level.

METHODS

Sample preparation and characterization. The CoO_x NP catalyst was prepared by the inverse micelle encapsulation route as described elsewhere.⁵² In short, CoCl_2 and poly(styrene-*b*-2-vinylpyridine) di-block copolymers were dissolved in toluene and the inverse micelles were dip-coated on an oxide-terminated Si(111) wafer. Subsequently, an O_2 plasma (20 W, 400 mTorr, 20 min) was applied to remove the polymer. Atomic force microscopy (Bruker MultiMode 8, tapping mode) was used to determine the NP height (7.0 ± 2.0 nm) and the average nearest neighbor distance (75 nm), see Figure 1 (b) and S1.

Photoemission measurements. The XPS measurements were carried out at the U49/2-PGM-1 beamline⁵³ located at the synchrotron radiation facility BESSY II in Berlin, Germany, using our NAP Sol³PES experimental setup.⁵⁴ The U49/2 beamline provides high-brilliant soft X-ray light up to a photon energy of $E_{\text{ph}} = 1400$ eV, enough to core-ionize water and also to excite Co 2p electrons into the valence band. The Sol³PES experimental setup is equipped with a Scienta Omicron R4000 HIP-2 hemispherical electron analyzer and is capable of detecting photoelectrons from samples inside our vacuum chamber at elevated pressures of up to 20 mbar water vapor. Since the detection screen and the multichannel plate of the electron-analyzer require high vacuum conditions ($< 10^{-8}$ mbar), there are three differential pumping sections, each consisting of a pinhole and two turbo molecular pumps, to separate the detector unit from the experimental chamber. A similar differential pumping scheme is used to deal with the pressure difference between the experimental chamber and the ultra-high vacuum ($< 10^{-9}$ mbar) inside the beamline. The distance between the electron analyzer entrance pinhole and the focal point of the soft X-ray beam is shorter than 500 μm , minimizing electron scattering at the water vapor in the chamber that causes an

attenuated electron signal. A detailed description of the SOL³PES setup and its characteristics can be found elsewhere.⁵⁴ During our measurements the angle between the polarization axis of the incoming soft X-ray and the electron analyzer was 54,7° (magic-angle geometry) to avoid any photoelectron angular distribution effects. The X-ray photoelectron spectra were energy calibrated using the Si 2p_{3/2} peak, with a 99.4 eV binding energy, and the spectra were normalized for comparison to the background intensity near 35 eV binding energy.⁵⁵

Resonant photoelectron spectroscopy. Our resonant photoelectron spectroscopy measurements focus on three specific Auger decay channels, 2p3p3p, 2p3p3d & 2p3d3d, and presented in Figure 1 (c).

In general, the excitation of a core electron into an unoccupied state (step 1 – black arrow), leads to a spectator Auger decay, where the hole is refilled by another energetically higher lying electron (step 2 – brown arrow) and the excess energy is used to kick out an electron of the sample (step 3 – grey arrow).^{34,54} A second resonant Auger channel is the participator Auger decay. Here, the first step is an excitation of a core-level electron into the valence band (2p→3d) similar to the spectator Auger decay. However, in the second step the excited electron itself refills the core hole and transfers its energy to another bound electron. As a consequence, the emitted Auger electron (step 3) will have the same kinetic energy as a directly emitted photoelectron, leading to a constructive interference of two identical final states enhancing the measured signals.^{34,46}

Partial electron yield x-ray absorption spectrum (PEY-XAS). Plotting the integral of a specific Auger decay channel signal as a function of the photon energy generates the respective (2p3p3p & 2p3p3d) partial electron yield x-ray absorption spectrum (PEY-XAS).^{32,34,56} These useful spectra exhibit signatures characteristic of the cobalt oxidation state and its coordination structure.^{12,15,30,31,41–44,57}

Near-ambient pressure. To create a 4-mbar water vapor environment we connected a test tube filled with Milli-Q water (degassed by two freeze-pump-thaw cycles) to the vacuum chamber. A needle valve is used to set the required H₂O gas flow into the chamber. To create a defined chamber pressure, the pumping flow rate is controlled with a motorized gas regulating valve (Pfeiffer EVR 116), which is managed by a proportional-integral-derivative controller (Pfeiffer RVC300 pressure gauge). Additionally, our sample holder consists of a brass block with pipes flushed with cold liquid ethanol, temperature-controlled by a chiller (Julabo F12-ED). The relative humidity in our chamber was calculated to be 70%, which translates into an adsorbed water layer thickness between 1.5 monolayer⁵⁵ and 5 monolayers⁵⁸.

ASSOCIATED CONTENT

Supporting Information. Atomic force microscopy image of the sample. NAP-XPS of SiO_x/Si(111) substrate. ‘On-resonant’ and ‘off-resonant’ spectra used to obtain the difference spectra in Figure 4 and additional O1s and C1s XP spectra.

The following file is available free of charge via the Internet at <http://pubs.acs.org>.

SI.pdf

AUTHOR INFORMATION

Corresponding Author

Email: Robert.seidel@helmholtz-berlin.de (RS)

Author Contributions

RS, AB, GW, DH performed the measurement at the BESSY II synchrotron. MB, AB and BRC prepared the samples. RS and DH wrote the manuscript, with comments from all authors.

ACKNOWLEDGMENT

DH, GW, and RS acknowledge funding from the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) through an Emmy-Noether grant (SE 2253/3-1). AB and BRC acknowledge financial support from DFG– project no. 388390466 – TRR 247, subproject A4. All authors thank the HZB staff for support during the beamtimes at BESSY II.

REFERENCES

- (1) Wang, J.; Cui, W.; Liu, Q.; Xing, Z.; Asiri, A. M.; Sun, X. Recent Progress in Cobalt-Based Heterogeneous Catalysts for Electrochemical Water Splitting. *Adv. Mater.* **2016**, *28*, 215–230.
- (2) Yan, Y.; Xia, B. Y.; Zhao, B.; Wang, X. A Review on Noble-Metal-Free Bifunctional Heterogeneous Catalysts for Overall Electrochemical Water Splitting. *J. Mater. Chem. A* **2016**, *4*, 17587–17603.
- (3) Wu, L.; Li, Q.; Wu, C. H.; Zhu, H.; Mendoza-Garcia, A.; Shen, B.; Guo, J.; Sun, S. Stable Cobalt Nanoparticles and Their Monolayer Array as an Efficient Electrocatalyst for Oxygen Evolution Reaction. *J. Am. Chem. Soc.* **2015**, *137*, 7071–7074.
- (4) Chou, N. H.; Ross, P. N.; Bell, A. T.; Tilley, T. D. Comparison of Cobalt-Based Nanoparticles as Electrocatalysts for Water Oxidation. *ChemSusChem* **2011**, *4*, 1566–1569.

- (5) Zhang, M.; De Respinis, M.; Frei, H. Time-Resolved Observations of Water Oxidation Intermediates on a Cobalt Oxide Nanoparticle Catalyst. *Nat. Chem.* **2014**, *6*, 362–367.
- (6) Zhang, Y.; Rosen, J.; Hutchings, G. S.; Jiao, F. Enhancing Photocatalytic Oxygen Evolution Activity of Cobalt-Based Spinel Nanoparticles. *Catal. Today* **2014**, *225*, 171–176.
- (7) Liao, L.; Zhang, Q.; Su, Z.; Zhao, Z.; Wang, Y.; Li, Y.; Lu, X.; Wei, D.; Feng, G.; Yu, Q.; Cai, X.; Zhao, J.; Ren, Z.; Fang, H.; Robles-Hernandez, F.; Baldelli, S.; Bao, J. Efficient Solar Water-Splitting Using a Nanocrystalline CoO Photocatalyst. *Nat. Nanotechnol.* **2014**, *9*, 69–73.
- (8) Jiao, F.; Frei, H. Nanostructured Cobalt and Manganese Oxide Clusters as Efficient Water Oxidation Catalysts. *Energy Environ. Sci.* **2010**, *3*, 1018–1027.
- (9) Esswein, A. J.; Mcmurdo, M. J.; Ross, P. N.; Bell, A. T.; Tilley, T. D. Size-Dependent Activity of Co₃O₄ Nanoparticle Anodes for Alkaline Water Electrolysis. *J. Phys. Chem. C* **2009**, *113*, 15068–15072.
- (10) Grzelczak, M.; Zhang, J.; Pfrommer, J.; Hartmann, J.; Driess, M.; Antonietti, M.; Wang, X. Electro- and Photochemical Water Oxidation on Ligand-Free Co₃O₄ Nanoparticles with Tunable Sizes. *ACS Catal.* **2013**, *3*, 383–388.
- (11) Zhuang, Z.; Sheng, W.; Yan, Y. Synthesis of Monodisperse Au@Co₃O₄ Core-Shell Nanocrystals and Their Enhanced Catalytic Activity for Oxygen Evolution Reaction. *Adv. Mater.* **2014**, *26*, 3950–3955.
- (12) Bergmann, A.; Jones, T. E.; Martinez Moreno, E.; Teschner, D.; Chernev, P.; Glich, M.; Reier, T.; Dau, H.; Strasser, P. Unified Structural Motifs of the Catalytically Active State of

Co(Oxyhydr)Oxides during the Electrochemical Oxygen Evolution Reaction. *Nat. Catal.* **2018**, *1*, 711–719.

(13) Chen, J.; Wu, X.; Selloni, A. Electronic Structure and Bonding Properties of Cobalt Oxide in the Spinel Structure. *Phys. Rev. B - Condens. Matter Mater. Phys.* **2011**, *83*, 1–7.

(14) Risbud, A. S.; Snedeker, L. P.; Elcombe, M. M.; Cheetham, A. K.; Seshadri, R. Wurtzite CoO. *Chem. Mater.* **2005**, *17*, 834–838.

(15) Papaefthimiou, V.; Dintzer, T.; Dupuis, V.; Tamion, A.; Tournus, F.; Hillion, A.; Teschner, D.; Hävecker, M.; Knop-Gericke, A.; Schlögl, R.; Zafeirotos, S. Nontrivial Redox Behavior of Nanosized Cobalt: New Insights from Ambient Pressure X-Ray Photoelectron and Absorption Spectroscopies. *ACS Nano* **2011**, *5*, 2182–2190.

(16) Wang, H. Y.; Hung, S. F.; Hsu, Y. Y.; Zhang, L.; Miao, J.; Chan, T. S.; Xiong, Q.; Liu, B. *In Situ* Spectroscopic Identification of μ -OO Bridging on Spinel Co₃O₄ Water Oxidation Electrocatalyst. *J. Phys. Chem. Lett.* **2016**, *7*, 4847–4853.

(17) Wang, H. Y.; Hung, S. F.; Chen, H. Y.; Chan, T. S.; Chen, H. M.; Liu, B. In Operando Identification of Geometrical-Site-Dependent Water Oxidation Activity of Spinel Co₃O₄. *J. Am. Chem. Soc.* **2016**, *138*, 36–39.

(18) Bajdich, M.; García-Mota, M.; Vojvodic, A.; Nørskov, J. K.; Bell, A. T. Theoretical Investigation of the Activity of Cobalt Oxides for the Electrochemical Oxidation of Water. *J. Am. Chem. Soc.* **2013**, *135*, 13521–13530.

- (19) Bergmann, A.; Martinez-Moreno, E.; Teschner, D.; Chernev, P.; Gliech, M.; De Araújo, J. F.; Reier, T.; Dau, H.; Strasser, P. Reversible Amorphization and the Catalytically Active State of Crystalline Co₃O₄ during Oxygen Evolution. *Nat. Commun.* **2015**, *6*, 1–9.
- (20) Tung, C. W.; Hsu, Y. Y.; Shen, Y. P.; Zheng, Y.; Chan, T. S.; Sheu, H. S.; Cheng, Y. C.; Chen, H. M. Reversible Adapting Layer Produces Robust Single-Crystal Electrocatalyst for Oxygen Evolution. *Nat. Commun.* **2015**, *6*, 1–9.
- (21) Favaro, M.; Yang, J.; Nappini, S.; Magnano, E.; Toma, F. M.; Crumlin, E. J.; Yano, J.; Sharp, I. D. Understanding the Oxygen Evolution Reaction Mechanism on CoO_x Using Operando Ambient-Pressure X-Ray Photoelectron Spectroscopy. *J. Am. Chem. Soc.* **2017**, *139*, 8960–8970.
- (22) Liu, J.; Bao, H.; Zhang, B.; Hua, Q.; Shang, M.; Wang, J.; Jiang, L. Geometric Occupancy and Oxidation State Requirements of Cations in Cobalt Oxides for Oxygen Reduction Reaction. *ACS Appl. Mater. Interfaces* **2019**, *11*, 12525–12534.
- (23) Grimes, R. W.; Lagerlof, K. P. D. Polymorphs of Cobalt Oxide. *J. Am. Ceram. Soc.* **1991**, *74*, 270–273.
- (24) Meyerheim, H. L.; Tusche, C.; Ernst, A.; Ostanin, S.; Maznichenko, I. V.; Mohseni, K.; Jedrecy, N.; Zegenhagen, J.; Roy, J.; Mertig, I.; Kirschner, J. Wurtzite-Type CoO Nanocrystals in Ultrathin ZnCoO Films. *Phys. Rev. Lett.* **2009**, *102*, 1–4.
- (25) Nam, K. M.; Shim, J. H.; Han, D. W.; Kwon, H. S.; Kang, Y. M.; Li, Y.; Song, H.; Seo, W. S.; Park, J. T. Syntheses and Characterization of Wurtzite CoO, Rocksalt CoO, and Spinel Co₃O₄ Nanocrystals: Their Interconversion and Tuning of Phase and Morphology. *Chem. Mater.* **2010**, *22*, 4446–4454.

- (26) Seo, W. S.; Shim, J. H.; Oh, S. J.; Lee, E. K.; Hur, N. H.; Park, J. T. Phase- and Size-Controlled Synthesis of Hexagonal and Cubic CoO Nanocrystals. *J. Am. Chem. Soc.* **2005**, *127*, 6188–6189.
- (27) Meyer, W.; Hock, D.; Biedermann, K.; Gubo, M.; Müller, S.; Hammer, L.; Heinz, K. Coexistence of Rocksalt and Wurtzite Structure in Nanosized CoO Films. *Phys. Rev. Lett.* **2008**, *101*, 1–4.
- (28) Meyer, W.; Biedermann, K.; Gubo, M.; Hammer, L.; Heinz, K. Superstructure in the Termination of CoO(111) Surfaces: Low-Energy Electron Diffraction and Scanning Tunneling Microscopy. *Phys. Rev. B - Condens. Matter Mater. Phys.* **2009**, *79*, 1–4.
- (29) Lukashuk, L.; Föttinger, K.; Kolar, E.; Rameshan, C.; Teschner, D.; Hävecker, M.; Knop-Gericke, A.; Yigit, N.; Li, H.; McDermott, E.; Stöger-Pollach, M.; Rupprechter, G. Operando XAS and NAP-XPS Studies of Preferential CO Oxidation on Co₃O₄ and CeO₂-Co₃O₄ Catalysts. *J. Catal.* **2016**, *344*, 1–15.
- (30) Papaefthimiou, V.; Dintzer, T.; Dupuis, V.; Tamion, A.; Tournus, F.; Teschner, D.; Hävecker, M.; Knop-Gericke, A.; Schlögl, R.; Zafeiratos, S. When a Metastable Oxide Stabilizes at the Nanoscale: Wurtzite CoO Formation upon Dealloying of PtCo Nanoparticles. *J. Phys. Chem. Lett.* **2011**, *2*, 900–904.
- (31) Turczyniak, S.; Luo, W.; Papaefthimiou, V.; Ramgir, N. S.; Haevecker, M.; MacHocki, A.; Zafeiratos, S. A Comparative Ambient Pressure X-Ray Photoelectron and Absorption Spectroscopy Study of Various Cobalt-Based Catalysts in Reactive Atmospheres. *Top. Catal.* **2016**, *59*, 532–542.

(32) Seidel, R.; Ghadimi, S.; Lange, K. M.; Bonhommeau, S.; Soldatov, M. A.; Golnak, R.; Kothe, A.; Könnecke, R.; Soldatov, A.; Thürmer, S.; Winter, B.; Aziz, E. F. Origin of Dark-Channel X-Ray Fluorescence from Transition-Metal Ions in Water. *J. Am. Chem. Soc.* **2012**, *134*, 1600–1605.

(33) Tian, L.; Zhu, J. L.; Chen, L.; An, B.; Liu, Q. Q.; Huang, K. L. Synthesis and Characterization of α -Cobalt Hydroxide Nanobelts. *J. Nanoparticle Res.* **2011**, *13*, 3483–3488.

(34) Lalithambika, S. S. N.; Atak, K.; Seidel, R.; Neubauer, A.; Brandenburg, T.; Xiao, J.; Winter, B.; Aziz, E. F. Chemical Bonding in Aqueous Hexacyano Cobaltate from Photon- and Electron-Detection Perspectives. *Sci. Rep.* **2017**, *7*, 1-13.

(35) Seidel, R.; Ghadimi, S.; Lange, K. M.; Bonhommeau, S.; Soldatov, M. A.; Golnak, R.; Kothe, A.; Könnecke, R.; Soldatov, A.; Thürmer, S.; Winter, B.; Aziz, E. F. Origin of Dark-Channel X-Ray Fluorescence from Transition-Metal Ions in Water. *J. Am. Chem. Soc.* **2012**, *134*, 1600–1605.

(36) Yamamoto, S.; Bluhm, H.; Andersson, K.; Ketteler, G.; Ogasawara, H.; Salmeron, M.; Nilsson, A. *In Situ* X-Ray Photoelectron Spectroscopy Studies of Water on Metals and Oxides at Ambient Conditions. *J. Phys. Condens. Matter* **2008**, *20*, 1–14.

(37) Kendelewicz, T.; Kaya, S.; Newberg, J. T.; Bluhm, H.; Mulakaluri, N.; Moritz, W.; Scheffler, M.; Nilsson, A.; Pentcheva, R.; Brown, G. E. Density Functional Theory Study on the Interaction of CO₂ with Fe₃O₄ (111) Surface. *Appl. Surf. Sci.* **2016**, *378*, 270–276.

(38) Yamamoto, S.; Kendelewicz, T.; Newberg, J. T.; Ketteler, G.; Starr, D. E.; Mysak, E. R.; Andersson, K. J.; Ogasawara, H.; Bluhm, H.; Salmeron, M.; Brown, G. E.; Nilsson, A. Water

Adsorption on α -Fe₂O₃ (0001) at near Ambient Conditions. *J. Phys. Chem. C* **2010**, *114*, 2256–2266.

(39) Uldry, A.; Vernay, F.; Delley, B. Systematic Computation of Crystal-Field Multiplets for x-Ray Core Spectroscopies. *Phys. Rev. B - Condens. Matter Mater. Phys.* **2012**, *85*, 1–14.

(40) Standard Reference Data, N. NIST Standard Reference Database 71. 1999. <https://doi.org/10.18434/T48C78>.

(41) Bazin, D.; Kovács, I.; Guzzi, L.; Parent, P.; Laffon, C.; De Groot, F.; Ducreux, O.; Lynch, J. Genesis of Co/SiO₂ Catalysts: XAS Study at the Cobalt LIII,II Absorption Edges. *J. Catal.* **2000**, *189*, 456–462.

(42) Hibberd, A. M.; Doan, H. Q.; Glass, E. N.; De Groot, F. M. F.; Hill, C. L.; Cuk, T. Co Polyoxometalates and a Co₃O₄ Thin Film Investigated by L-Edge X-Ray Absorption Spectroscopy. *J. Phys. Chem. C* **2015**, *119*, 4173–4179.

(43) Morales, F.; De Groot, F. M. F.; Glatzel, P.; Kleimenov, E.; Bluhm, H.; Hävecker, M.; Knop-Gericke, A.; Weckhuysen, B. M. *In Situ* X-Ray Absorption of Co/Mn/TiO₂ Catalysts for Fischer-Tropsch Synthesis. *J. Phys. Chem. B* **2004**, *108*, 16201–16207.

(44) Zafeiratos, S.; Dintzer, T.; Teschner, D.; Blume, R.; Hävecker, M.; Knop-Gericke, A.; Schlögl, R. Methanol Oxidation over Model Cobalt Catalysts: Influence of the Cobalt Oxidation State on the Reactivity. *J. Catal.* **2010**, *269*, 309–317.

(45) Ueda, Y.; Sato, H.; Hirai, C.; Arita, M.; Kimura, A.; Taniguchi, M.; Irisawa, A.; Muro, T.; Saitoh, Y.; Sekiyama, A.; Suga, S. Co 2p-3d Resonant Photoemission Spectroscopy of CoSb₃. *J. Electron Spectros. Relat. Phenom.* **2005**, *144–147*, 663–666.

- (46) Tanaka, A.; Jo, T. Resonant 3 d , 3 p and 3 s Photoemission in Transition Metal Oxides Predicted at 2 p Threshold. *J. Phys. Soc. Japan* **1994**, *63*, 2788–2807.
- (47) Winter, B.; Weber, R.; Widdra, W.; Dittmar, M.; Faubel, M.; Hertel, I. V. Full Valence Band Photoemission from Liquid Water Using EUV Synchrotron Radiation. *J. Phys. Chem. A* **2004**, *108*, 2625–2632.
- (48) Jugnet, Y.; Duc, T. M. Structure Electronique Des Oxydes de Cobalt CoO et Co₃O₄. *J. Phys. Chem. Solids* **1979**, *40*, 29–37.
- (49) Alay, J. L.; Hirose, M. The Valence Band Alignment at Ultrathin SiO₂/Si Interfaces. *J. Appl. Phys.* **1997**, *81*, 1606–1608.
- (50) Velasco-Velez, J. J.; Pfeifer, V.; Hävecker, M.; Weatherup, R. S.; Arrigo, R.; Chuang, C. H.; Stotz, E.; Weinberg, G.; Salmeron, M.; Schlögl, R.; Knop-Gericke, A. Photoelectron Spectroscopy at the Graphene-Liquid Interface Reveals the Electronic Structure of an Electrodeposited Cobalt/Graphene Electrocatalyst. *Angew. Chem. Int. Ed.* **2015**, *54*, 14554–14558.
- (51) Velasco-Vélez, J. J.; Jones, T. E.; Streibel, V.; Hävecker, M.; Chuang, C. H.; Frevel, L.; Plodinec, M.; Centeno, A.; Zurutuza, A.; Wang, R.; Arrigo, R.; Mom, R.; Hofmann, S.; Schlögl, R.; Knop-Gericke, A. Electrochemically Active Ir NPs on Graphene for OER in Acidic Aqueous Electrolyte Investigated by *in Situ* and *ex Situ* Spectroscopies. *Surf. Sci.* **2019**, *681*, 1–8.
- (52) Bernal, M.; Bagger, A.; Scholten, F.; Sinev, I.; Bergmann, A.; Ahmadi, M.; Rossmeisl, J.; Cuenya, B. R. CO₂ Electroreduction on Copper-Cobalt Nanoparticles: Size and Composition Effect. *Nano Energy* **2018**, *53*, 27–36.

(53) Sawhney, K. J. S.; Senf, F.; Gudat, W. PGM Beamline with Constant Energy Resolution Mode for U49-2 Undulator at BESSY-II. *Nucl. Instrum. Methods Phys. Res., Sect. A* **2001**, *467–468*, 466–469.

(54) Seidel, R.; Pohl, M. N.; Ali, H.; Winter, B.; Aziz, E. F. Advances in Liquid Phase Soft-x-Ray Photoemission Spectroscopy: A New Experimental Setup at BESSY II. *Rev. Sci. Instrum.* **2017**, *88*, 1–8.

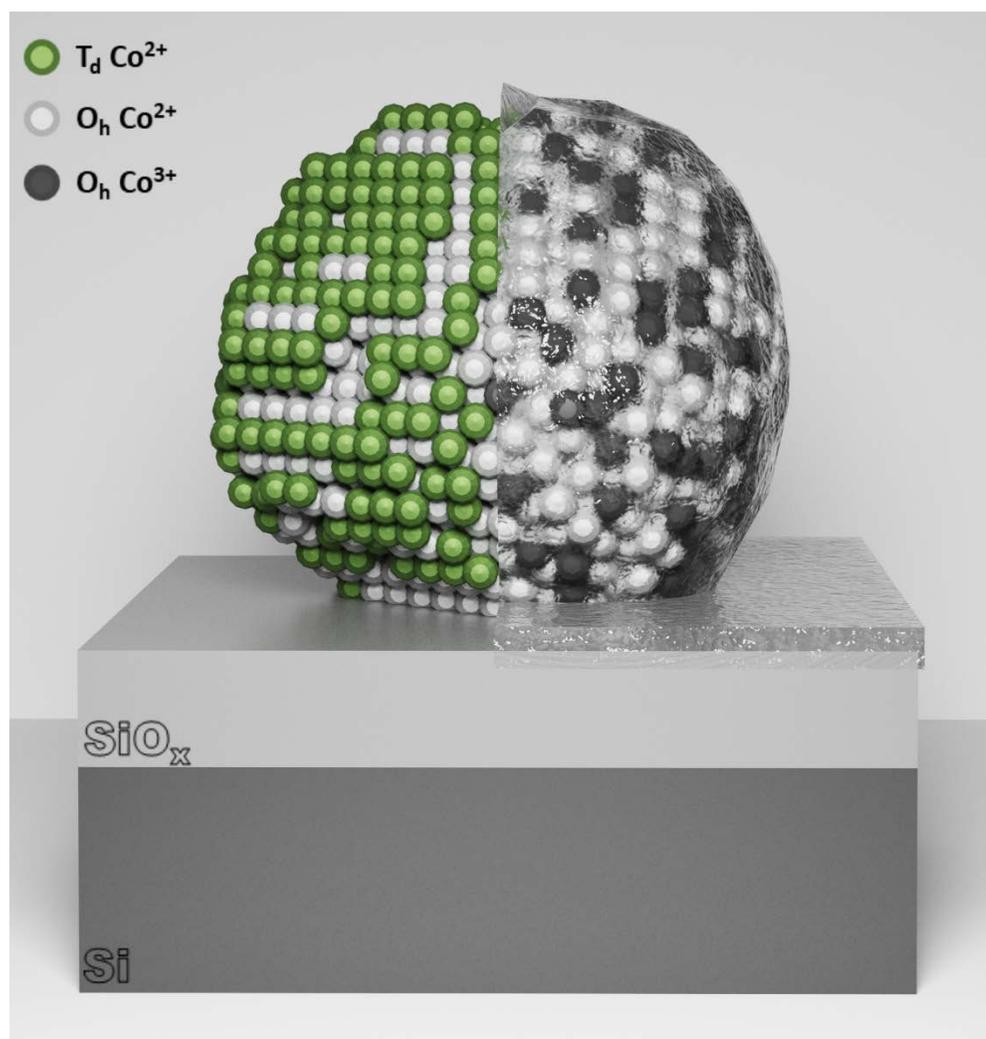
(55) Verdaguer, A.; Weis, C.; Oncins, G.; Ketteler, G.; Bluhm, H.; Salmero, M. Growth and Structure of Water on SiO₂ Films on Si Investigated by Kelvin Probe Microscopy and *in Situ* X-Ray Spectroscopies. *Langmuir* **2007**, *23*, 9699–9703.

(56) Soldatov, M. A.; Lange, K. M.; Gotz, M. D.; Engel, N.; Golnak, R.; Kothe, A.; Aziz, E. F. On the Origin of Dips in Total Fluorescence Yield X-Ray Absorption Spectra: Partial and Inverse Partial Fluorescence Yield at the L-Edge of Cobalt Aqueous Solution. *Chem. Phys. Lett.* **2012**, *546*, 164–167.

(57) Herranz, T.; Deng, X.; Cabot, A.; Guo, J.; Salmeron, M. Influence of the Cobalt Particle Size in the CO Hydrogenation Reaction Studied by *in Situ* X-Ray Absorption Spectroscopy. *J. Phys. Chem. B.* **2009**, *113*, 10721–10727.

(58) Ketteler, G.; Yamamoto, S.; Bluhm, H.; Andersson, K.; Starr, D. E.; Ogletree, D. F.; Ogasawara, H.; Nilsson, A.; Salmeron, M. The Nature of Water Nucleation Sites on TiO₂(110) Surfaces Revealed by Ambient Pressure X-Ray Photoelectron Spectroscopy. *J. Phys. Chem. C* **2007**, *111*, 8278–8282.

TABLE OF CONTENTS



Schematic comparison of **one** CoO_x NP on SiO_x/Si(111) substrate without (left) and with adsorbed H₂O (right).